

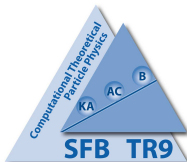
Moments of the Non-Diagonal Current Correlators

in QCD up to Three Loop Order

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Cosmology meets Particle Physics – Ideas & Measurements



Motivation

for this Calculation

- current correlators: building blocks

$m_1 = m_2$ heavy-heavy or diagonal

$\sigma(e^+e^- \rightarrow \text{had.}), \Gamma(Z \rightarrow \text{had.}), \Gamma(H \rightarrow t\bar{t})$

$m_1 \neq m_2$ heavy-light or non-diagonal

$\sigma(q\bar{q} \rightarrow t\bar{b}), \Gamma(W^+ \rightarrow t\bar{b}), \Gamma(H^+ \rightarrow q\bar{q})$

- moments: expansion coefficients in small q^2 ;
used for extraction of parameters

- $m_{c,b}$: diagonal vector correlator (at 4L)

\leftrightarrow experiment $\sigma(e^+e^- \rightarrow \text{hadrons})$

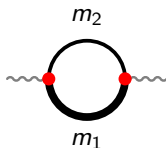
[Chetyrkin, Kühn, Maier, Maierhöfer, Marquard, Sturm, '07-'10]

- $m_{c,b}, \alpha_s$: diagonal pseudo-scalar correlator (at 4L)

\leftrightarrow moments from lattice simulations

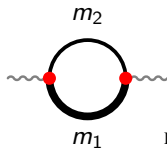
[Allison et al., HPQCD Collaboration, '08],

[McNeile et al., HPQCD Collaboration, '10]



Motivation

for this Calculation



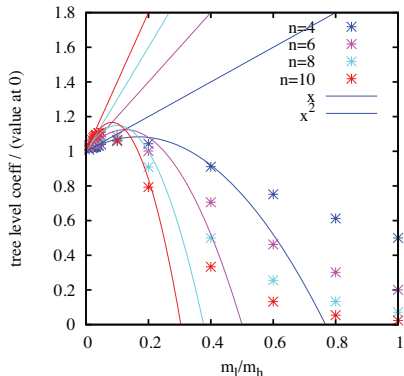
$$x = m_1/m_2$$

Idea

extract decay constants f_D and f_B from lattice data; study semileptonic processes [Koponen et al., HPQCD Collaboration, '10]

- zero-mass case **not sufficient**: m_c/m_b -corrections important
- lattice simulations for variety of masses (not only the physical ones)
- extrapolation to values of physical interest

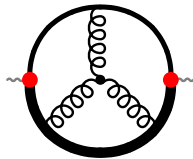
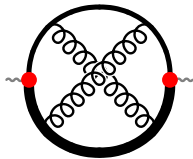
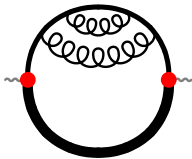
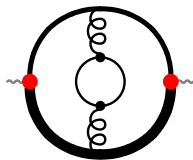
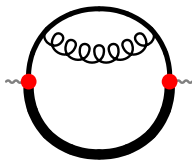
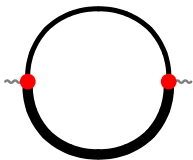
HL moment coeffs - exact vs expansion in m_l/m_h



Needed moments of h.-l. corr. with dependence on arbitrary masses

Sample Diagrams

at 1-, 2- and 3-Loop Order



(1, 3 and 34 diagrams)

Some Definitions

Notations & Properties

- scalar (s), pseudo-scalar (p), vector (v) and axial-vector (a) quark-currents:

$$\begin{aligned} j^s &= \bar{\psi}_1 \psi_2 & j_\mu^v &= \bar{\psi}_1 \gamma_\mu \psi_2 \\ j^p &= \bar{\psi}_1 i \gamma^5 \psi_2 & j_\mu^a &= \bar{\psi}_1 \gamma_\mu \gamma^5 \psi_2 \end{aligned}$$

- current correlators (polarization functions):

$$q^2 \Pi^\delta(q^2) = i \int dx e^{iqx} \langle 0 | T j^\delta(x) j^{\delta\dagger}(0) | 0 \rangle \quad (\delta = s, p)$$

$$\begin{aligned} (-q^2 g_{\mu\nu} + q_\mu q_\nu) \Pi_T^\delta(q^2) + q_\nu q_\mu \Pi_L^\delta(q^2) = \\ i \int dx e^{iqx} \langle 0 | T j_\mu^\delta(x) j_\nu^{\delta\dagger}(0) | 0 \rangle \quad (\delta = v, a) \end{aligned}$$

Some Definitions

Notations & Properties

- dimensionless variables: $z = q^2/m_1^2$, $x = m_2/m_1$

$$\bar{\Pi}^\delta(q^2) = \frac{3}{16\pi^2} \sum_{n \geq -1} \bar{C}_n^\delta(x) z^n$$

- perturbative series:

$$\bar{C}_n^\delta = \bar{C}_n^{(0),\delta} + \frac{\alpha_s}{\pi} \bar{C}_n^{(1),\delta} + \left(\frac{\alpha_s}{\pi}\right)^2 \underline{\bar{C}_n^{(2),\delta}} + \mathcal{O}(\alpha_s^3)$$

Some Definitions

Notations & Properties

Checks

- $m_2 = 0$:

$$\bar{\Pi}^s = \bar{\Pi}^p \quad \text{and} \quad \bar{\Pi}_{T,L}^v = \bar{\Pi}_{T,L}^a$$

- mirror symmetry (due to γ_5):

$$\bar{C}_n^s(x) = \bar{C}_n^p(-x) \quad \text{and} \quad \bar{C}_n^v(x) = \bar{C}_n^a(-x)$$

- Ward-identity ($x = m_2/m_1$):

$$\bar{C}_{L,n}^v = (1-x)^2 \bar{C}_{n+1}^s \quad \text{and} \quad \bar{C}_{L,n}^a = (1+x)^2 \bar{C}_{n+1}^p$$

- at 2 loops: agreement with exact result [Djoudi, Gambino, '94/'95]
- at 3 loops: agreement with special cases $m_2 = m_1$ or $m_2 = 0$
[Chetyrkin, Kühn, Steinhauser, '97], [Chetyrkin, Steinhauser, '01]

Renormalization

of the Correlators

- 1 parameter-renormalization:

$$m_{1,2}^0 \rightarrow Z_m m_{1,2}$$

$$\alpha_s^0 \rightarrow Z_\alpha \alpha_s$$

- 2 renormalization of the current:

$$(j^\delta)_{\text{bare}} = Z_{j^\delta} (j^\delta)_{\text{ren.}} \quad \text{with} \quad \boxed{Z_{j^s,p} = Z_m} \quad \text{and} \quad \boxed{Z_{j^{v,a}} = 1}$$

- 3 QED-like overall-renormalization:

$$\boxed{\Pi^\delta(0) = 0} \quad \Leftarrow \quad \Pi^\delta(q^2) = \bar{\Pi}^\delta(q^2) - \frac{3}{16\pi^2} \left(\bar{C}_0^\delta + \frac{\bar{C}_{-1}^\delta}{z} \right)$$

Details of the Calculation

of two different Expansions

- highly automated setup: QGRAF, q2e, exp and MATAD
- 3 scales \Rightarrow two different expansions at 3 loops:
 - around $x = 0$: asymptotic expansion in $q^2 \ll m_2^2 \ll m_1^2$
 - around $x = 1$: naive expansion in q^2 and $\Delta m = m_2 - m_1$

$$\frac{1}{m_2 - \not{p}} = \sum_{i \geq 0} \frac{(-\Delta_m)^i}{(m_1 - \not{p})^{i+1}}$$

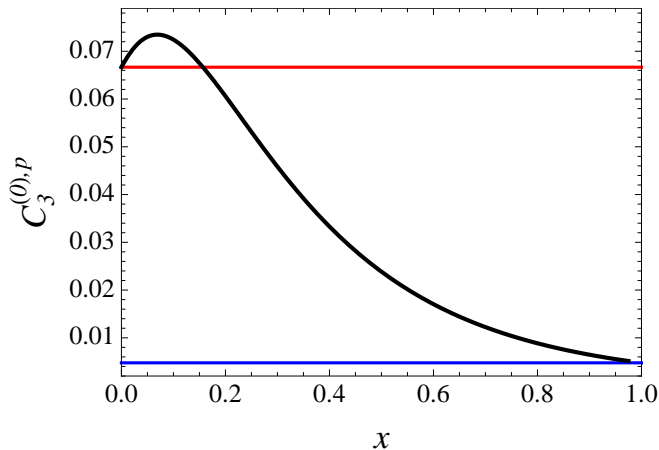
- RGE-methods: exact μ -dependence at 3 loops

$$\mu^2 \frac{d^2}{d\mu^2} \left(Z_{j^\delta}^2 \Pi^\delta \right) = 0$$

with $\Pi_{(2)}^\delta = \dots + l_\mu \dots + l_\mu^2 \dots$ and $l_\mu = \log \frac{\mu^2}{m_1^2}$

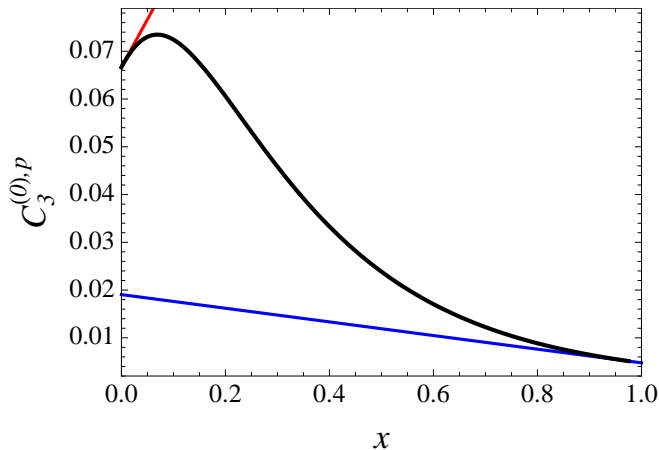
Results

at 1-Loop Order



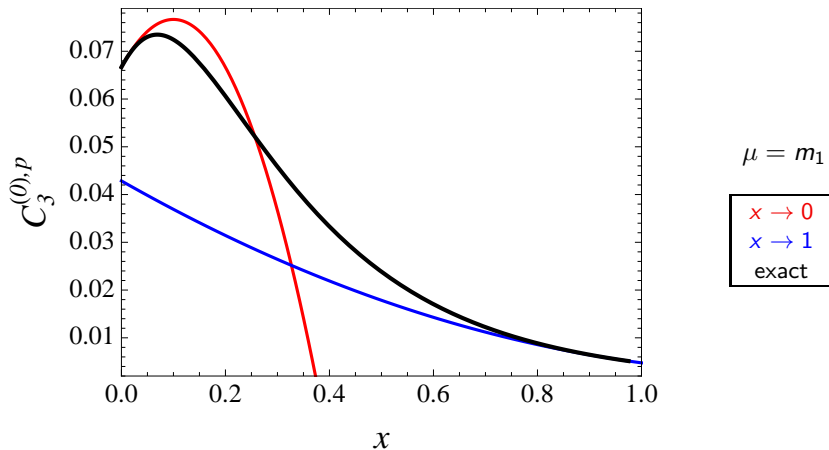
Results

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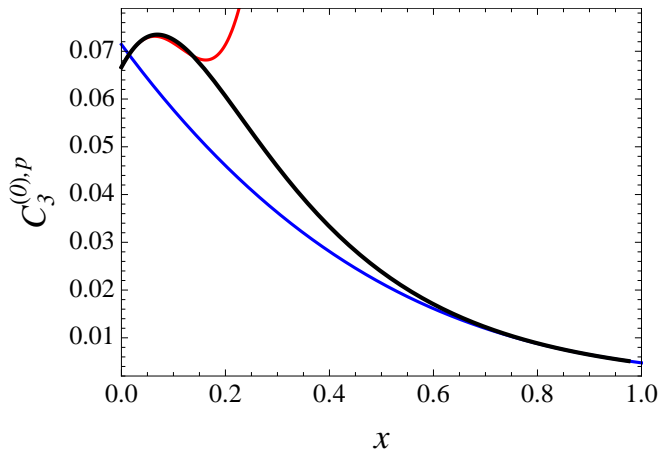
Results

at 1-Loop Order



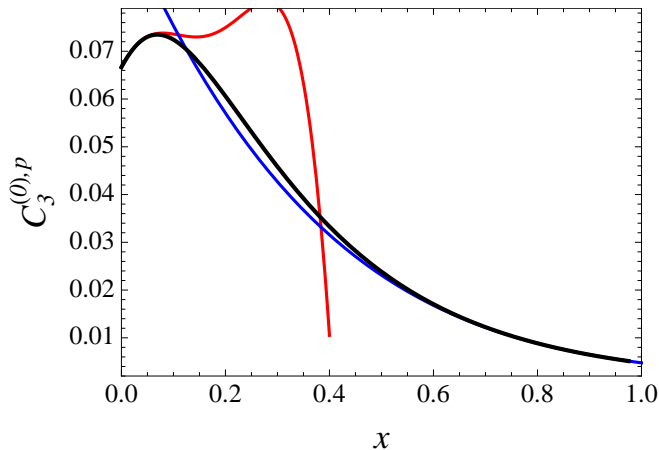
Results

at 1-Loop Order



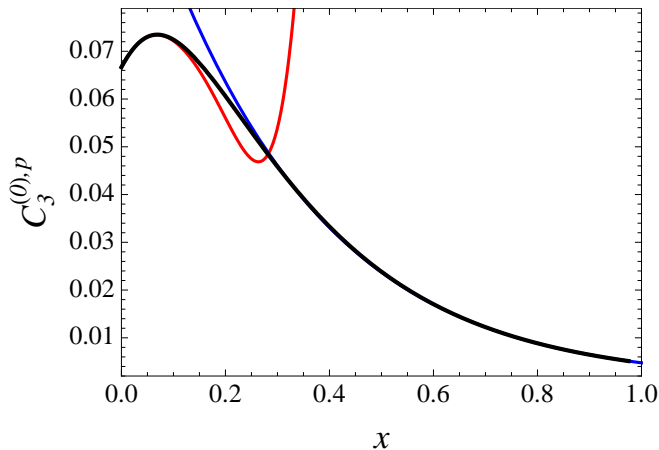
Results

at 1-Loop Order



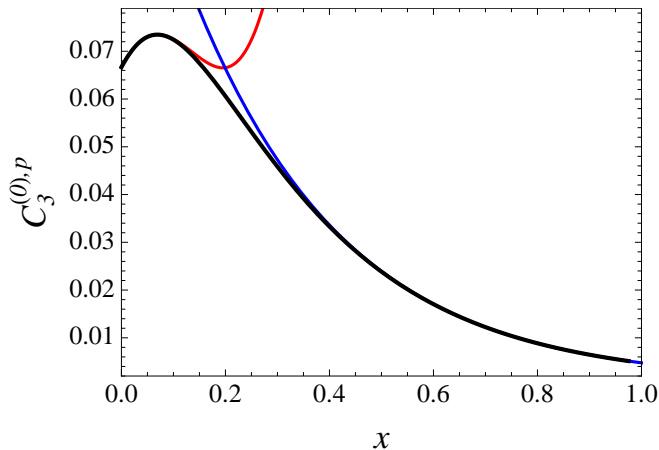
Results

at 1-Loop Order

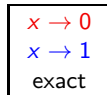


Results

at 1-Loop Order

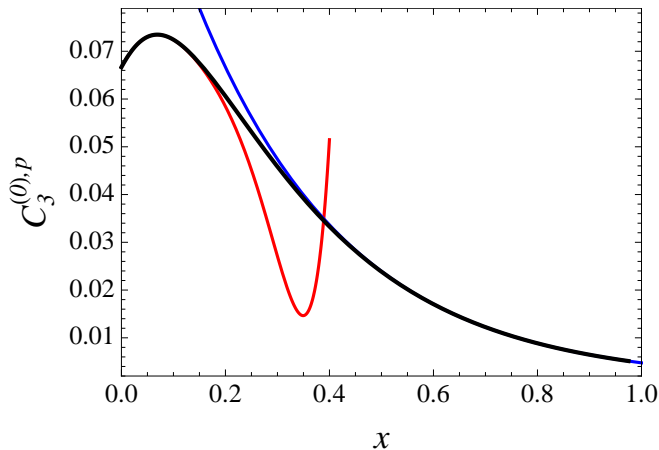


$\mu = m_1$



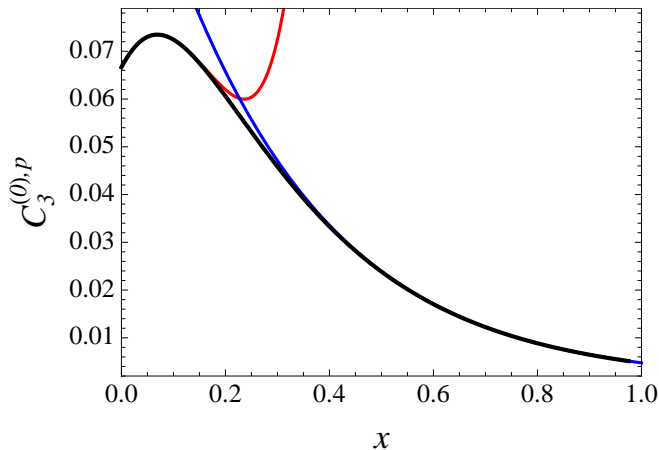
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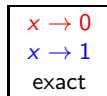


Results

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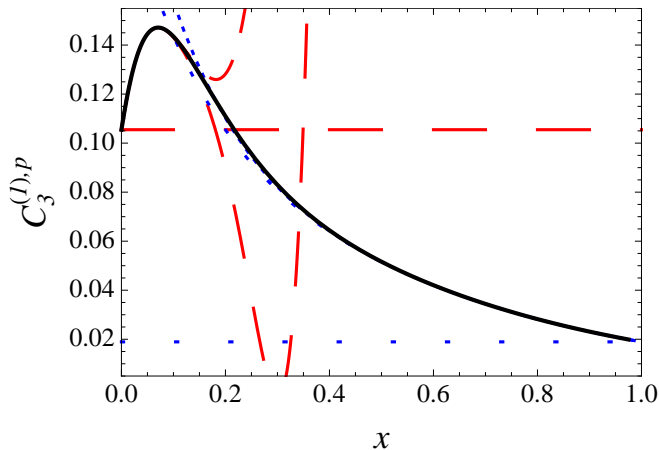


$\mu = m_1$



Results

at 2-Loop Order



$\mu = m_1$

$x \rightarrow 0$
$x \rightarrow 1$
exact

$\mathcal{O}(x^8), \mathcal{O}(x^9)$

$\mathcal{O}((1-x)^8),$

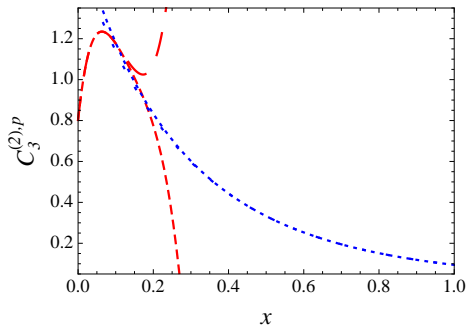
$\mathcal{O}((1-x)^9)$

Results

at 3-Loop Order

Method

- 1 find x_0 and x_1
(0.1, 0.5)
- 2 construct
interpolation
function
- 3 fit to polynomial
in $x^{1/2}$ up to
 $\mathcal{O}(x^3)$

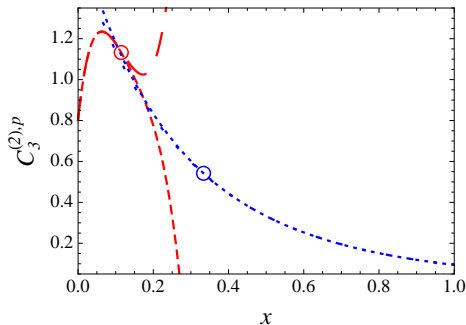


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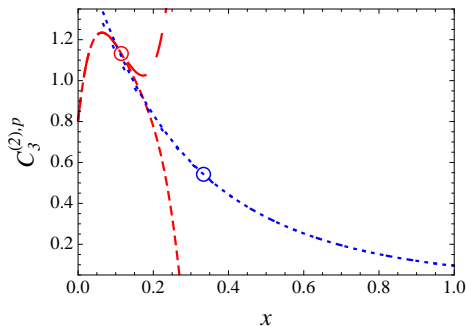


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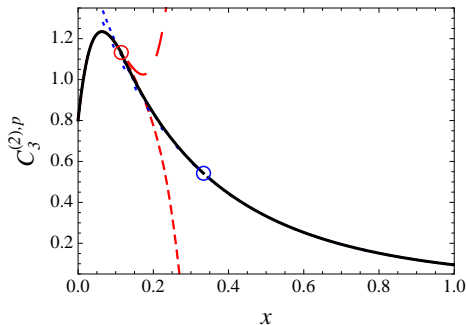


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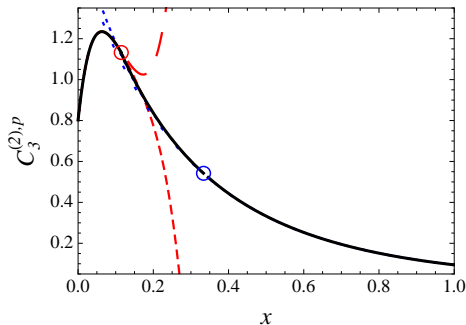


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Observation

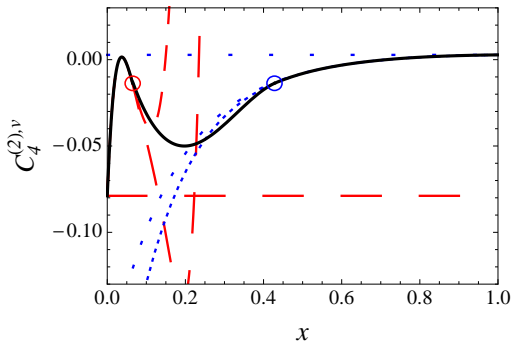
works perfectly at 1 and 2 loops \Rightarrow applied at 3 loop level

Results

at 3-Loop Order

only 2 cases (out of $4 \times 4 = 16$) are more difficult to approximate:

$C_3^{(2),v}$ and $C_4^{(2),v}$ (more compl. behaviour)

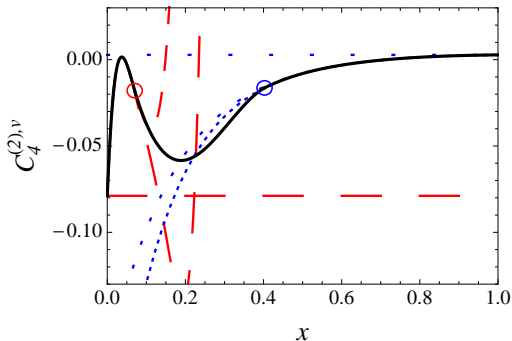


Results

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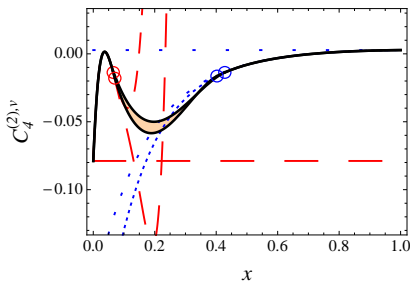


Results

at 3-Loop Order

only 2 cases (out of $4 \times 4 = 16$) are more difficult to approximate:

$C_3^{(2),v}$ and $C_4^{(2),v}$ (more compl. behaviour)



But our approach works fine in most cases
Error below 5% (estimated at 1 & 2 loops)

Summary & Outlook

Summary

- 4 cases (s, p, v, a); up to 4th moment
- 1 loop & 2 loops: exact x -dependence ($x = m_2/m_1$)
- 3 loops:
 - expansions around $x = 0$ and $x = 1$
 - combinations valid for arbitrary x
 - exact μ -dependence

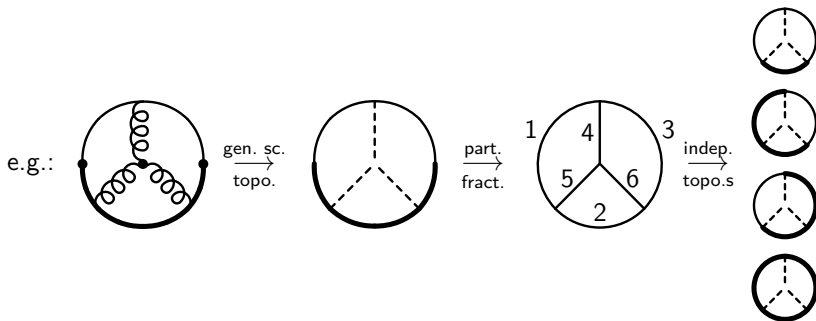
⇒ Mathematica package: `coefhl.m`

Outlook

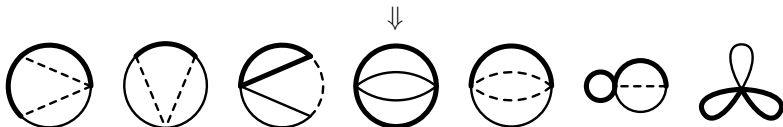
- ⇒ calculation retaining the exact x -dependence at 3 loops
- ⇒ go to higher moments

Summary & Outlook

Exact Calculation



Laporta reduction to master integrals



(\Rightarrow apply Method of Differential Equations)

- 1st order diff. eqn.s (in gen.: **chained**): e.g.

$$\frac{d}{dx} \begin{array}{c} x \\ \circlearrowleft \\ 0 \\ \circlearrowright \\ 1 \end{array} \sim \begin{array}{c} \bullet \\ \circlearrowleft \\ \circlearrowright \end{array} \xrightarrow[\text{red.}]{\text{Lap.}} \underbrace{\begin{array}{c} \circlearrowleft \\ \circlearrowright \end{array}}_{\text{hom.}}, \quad \underbrace{\infty}_{\text{inhom.}}$$

$$\frac{d}{dx} f(1, 1, 1) \sim f(1, 2, 1) \longrightarrow f(1, 1, 1), \quad f(1, 1, 0)$$

- expand in ϵ** (decouple) \Rightarrow solve by **Euler's method** using (known) 1-scale integrals as bounds; in this example

$$x \rightarrow 0: \begin{array}{c} \circlearrowleft \\ \text{---} \\ \circlearrowright \end{array} \quad \text{or} \quad x \rightarrow 1: \begin{array}{c} \circlearrowleft \\ \text{---} \\ \circlearrowright \end{array}$$

- use **Harmonic** Polylogs with: $f_1(x) = \frac{1}{1-x}$, $f_2(x) = \frac{1}{x}$, $f_3(x) = \frac{1}{1+x}$,

$$H(i, i_1, \dots, i_n; x) = \int_0^x f_i(y) H(i_1, \dots, i_n) dy$$